
Predictive freeze-in

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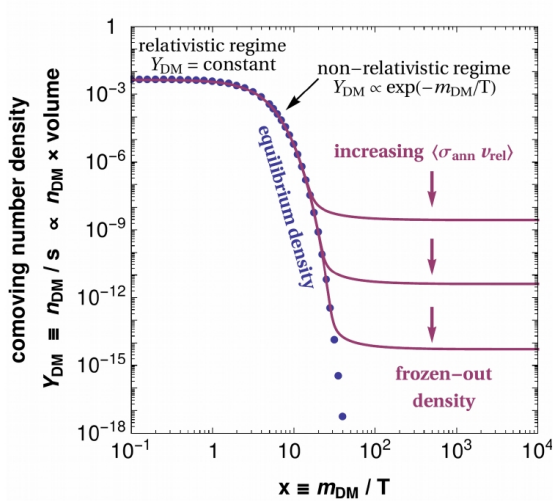
University of Helsinki

- gravitational particle production background
- freeze-in with low temperatures
- simplest models (Higgs portal, Z' , ...)
- direct DM detection, invisible Higgs decay

Dark Matter Models

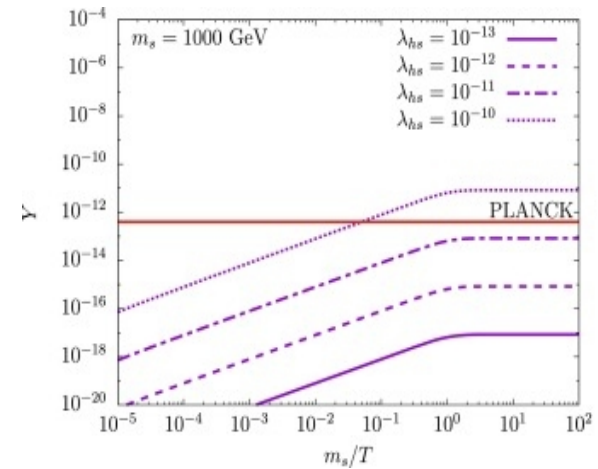
thermal

non-thermal



No memory

(“attractor solution”)



Memory

General remarks

– *psychologically thermal particles are natural:*

we observe ONLY thermal particles in reality (e, gamma, ...)

because we only see particles with gauge interactions

– *freeze-out is real (neutrinos)*

– *non-thermal particles ~ paradigm shift, challenging:*

- **initial conditions are as important as the production mechanism**

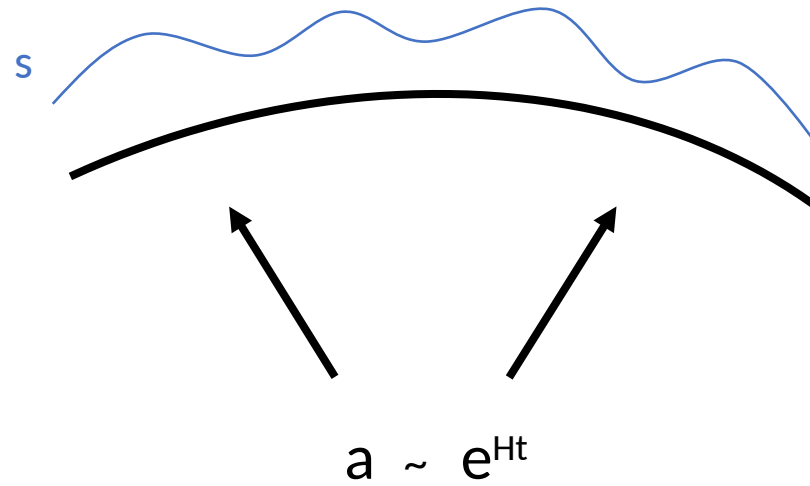
(or prove otherwise)

- **gravity is always there** → must prove it's irrelevant

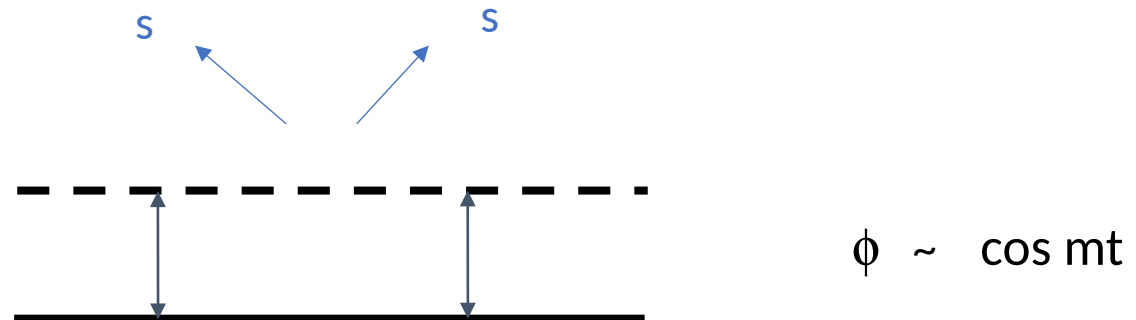
(otherwise there's nothing to talk about)

Gravitational particle production

Inflation:



After inflation:

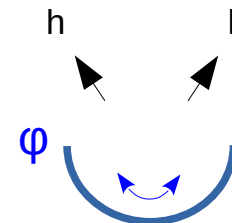


Example: perturbative calculation (no collective effects)

Take

$$\Delta V = \frac{1}{4} \lambda_{\phi h} \phi^2 h^2$$

$$\phi^2(t) = \sum_{n=-\infty}^{\infty} \zeta_n e^{-in\omega t}$$



Transition amplitude $|0\rangle \rightarrow |p, q\rangle$

$$-i \int_{-\infty}^{\infty} dt \langle f | V(t) | i \rangle = -i \frac{\lambda_{\phi h}}{2} (2\pi)^4 \delta(\mathbf{p} + \mathbf{q}) \sum_{n=1}^{\infty} \zeta_n \delta(E_p + E_q - n\omega)$$

Reaction rate per unit volume

$$\Gamma = \sum_{n=1}^{\infty} \Gamma_n = \sum_{n=1}^{\infty} \frac{1}{2} \int |\mathcal{M}_n|^2 d\Pi_n = \frac{\lambda_{\phi h}^2}{64\pi} \sum_{n=1}^{\infty} |\zeta_n|^2 \sqrt{1 - \left(\frac{2m_h}{n\omega}\right)^2} \theta(n\omega - 2m_h)$$

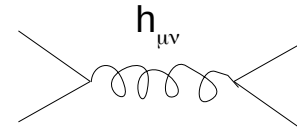
Same applies to Planck-suppressed op-s, e.g.

$$\frac{1}{M_{\text{Pl}}^2} \phi^4 s^2$$

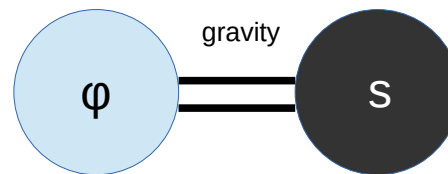
Note:

gravity is not classical at high energies

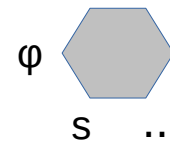
one-graviton exchange *inadequate*



Need



E.g. n-point function in strings



Planck-suppressed operators are very efficient in particle production:

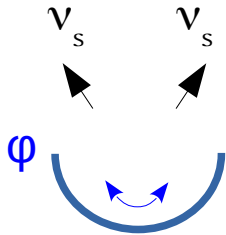
$$\frac{\phi^4 s^2}{M_{\text{Pl}}^2}, \quad \frac{\phi^6 s^2}{M_{\text{Pl}}^4}, \quad \frac{\phi^8 s^2}{M_{\text{Pl}}^6}, \quad \dots$$

coefficients unknown! (quantum gravity)

Fermions :

Koutroulis, OL, Pokorski '24

$$\frac{c}{M_{\text{Pl}}} \phi^2 \bar{\Psi} \Psi, \dots$$



*produces viable **COLD keV sterile neutrino DM** ($C \sim 0.1$)*

Freeze-in models suffer from the gravitational background → not predictive

- Gravitationally produced relics may be the end of the story
- If not, can get rid of it:

inflaton energy density $\sim a^{-3}$

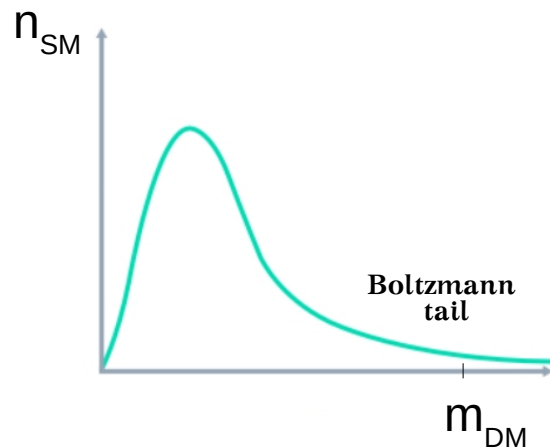
rel. relic energy density $\sim a^{-4}$



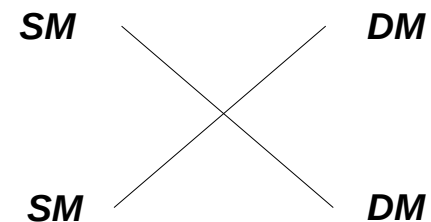
low T_R

What if $T_R < m_{DM}$?

Cosme, Costa, OL '23



off the
tail



Logic:

All we know

$$T_R > 4 \text{ MeV}$$

(*Hannestad 2004*)

High T_R = strong assumption !

Relax it and reanalyze

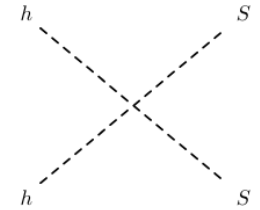
- Higgs portal
- Z' extension
- ν_R + scalar extension



direct DM detection + invisible Higgs decay

Simplest model = Higgs portal DM

$$V(s) = \frac{1}{2} \lambda_{hs} s^2 H^\dagger H + \frac{1}{2} m_s^2 s^2$$



Boltzmann equation:

$$\dot{n} + 3Hn = \Gamma(h_i h_i \rightarrow ss) - \Gamma(ss \rightarrow h_i h_i)$$

No annihilation:

$$\Gamma(h_i h_i \rightarrow ss) \simeq \frac{\lambda_{hs}^2 T^3 m_s}{27 \pi^4} e^{-2m_s/T} \quad \rightarrow \quad \lambda_{hs} \simeq 3 \times 10^{-11} e^{m_s/T_R} \sqrt{\frac{T_R}{m_s}}$$

With annihilation:

$$\Gamma(ss \rightarrow h_i h_i) = \sigma(ss \rightarrow h_i h_i) v_r n^2, \quad \sigma(ss \rightarrow h_i h_i) v_r = 4 \times \frac{\lambda_{hs}^2}{64 \pi m_s^2}$$

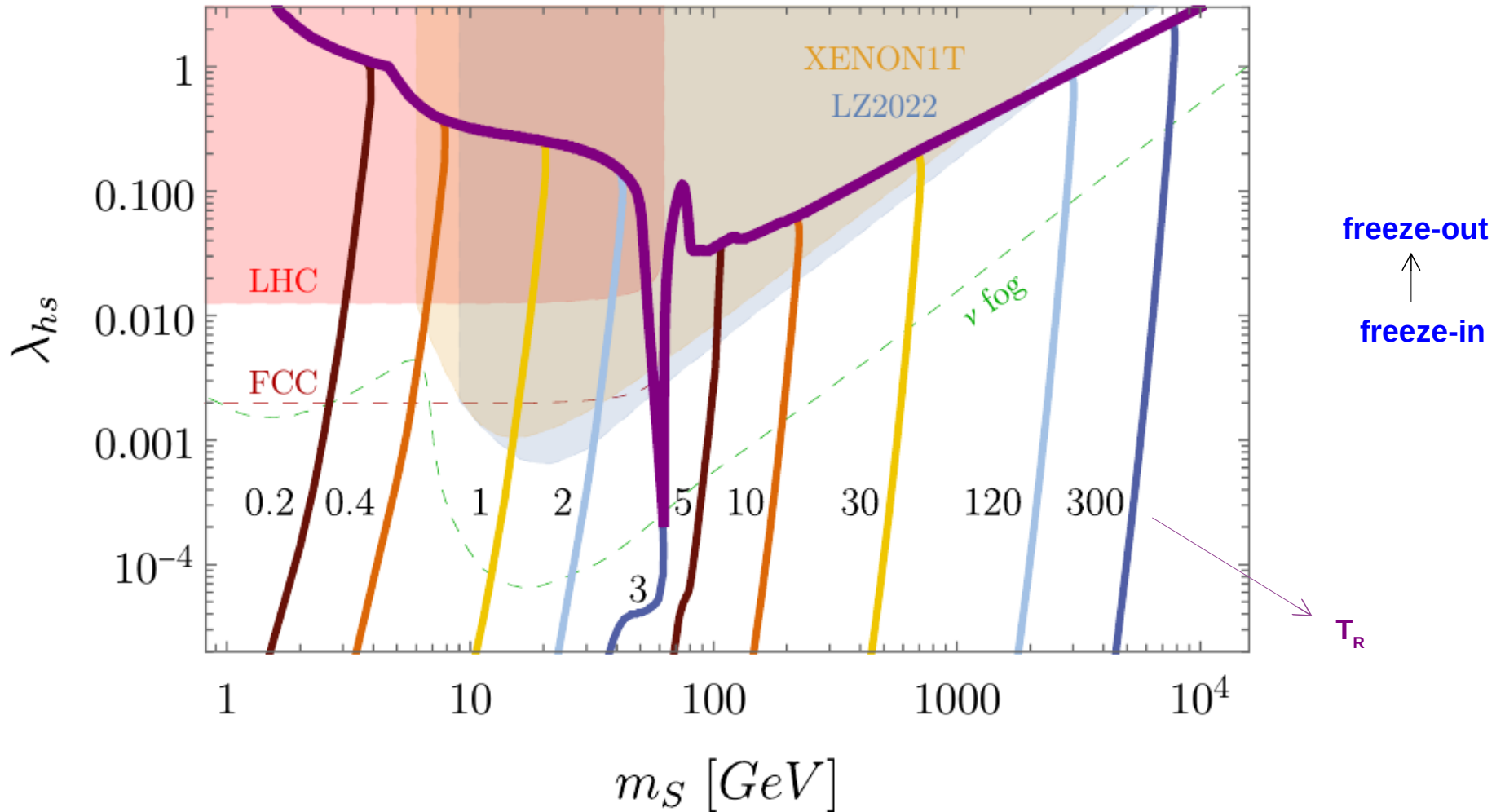
No thermalization:

$$\Gamma(h_i h_i \rightarrow ss) \neq \Gamma(ss \rightarrow h_i h_i)$$

Scalar Higgs portal DM :

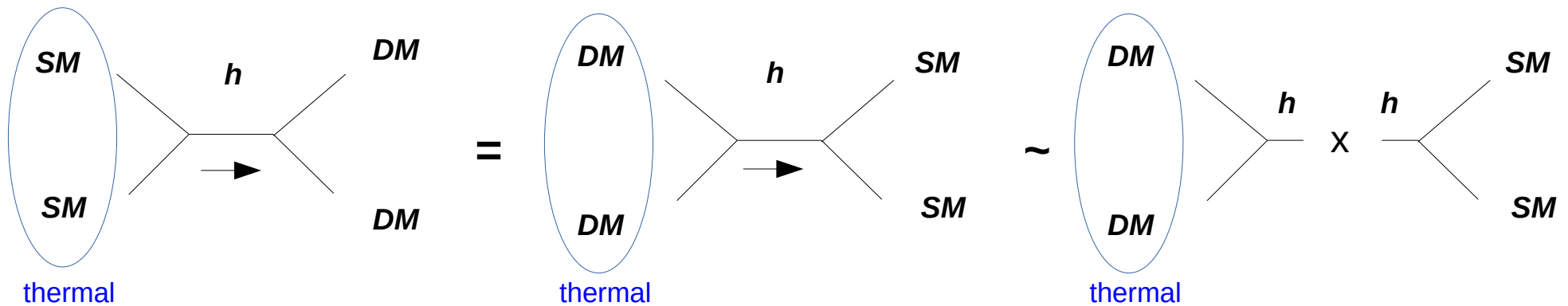
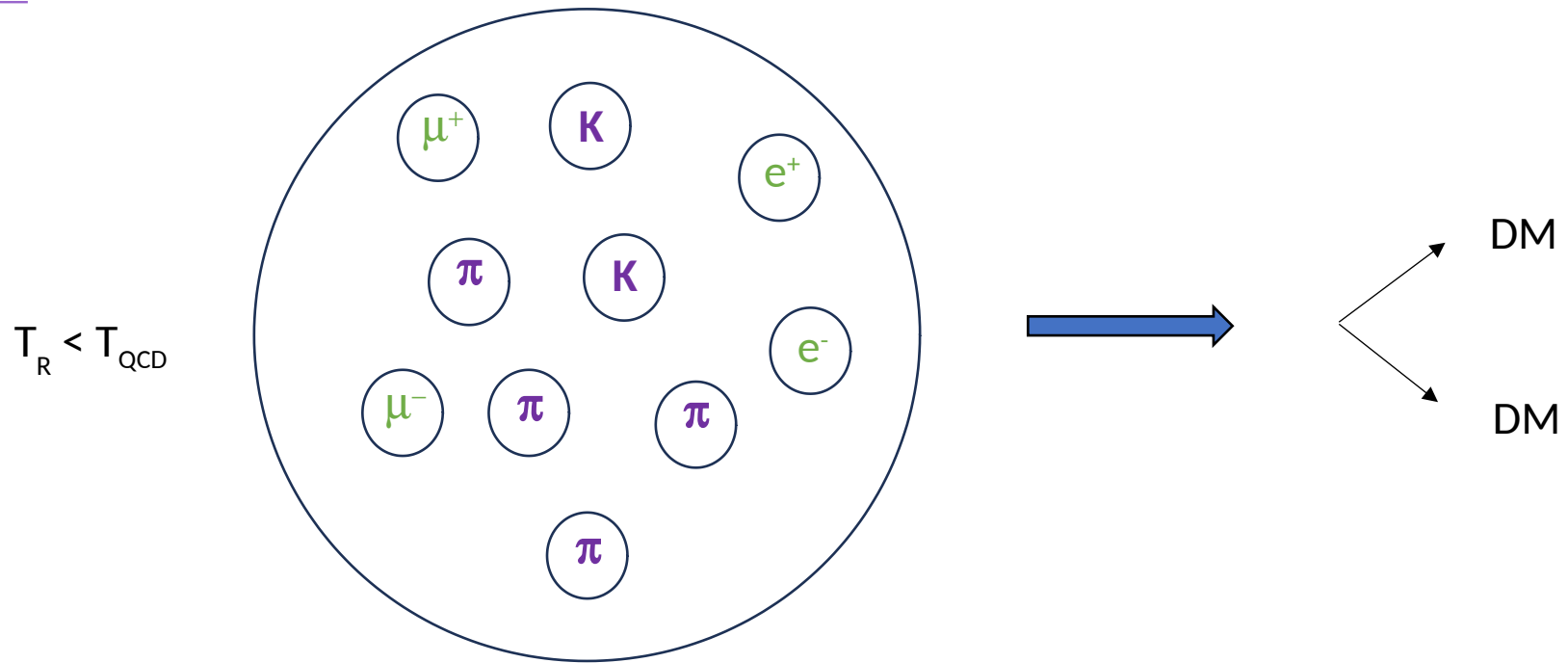
$$V(s) = \frac{1}{2}\lambda_{hs}s^2H^\dagger H + \frac{1}{2}m_s^2s^2$$

Arcadi, Costa, Goudelis, OL '24

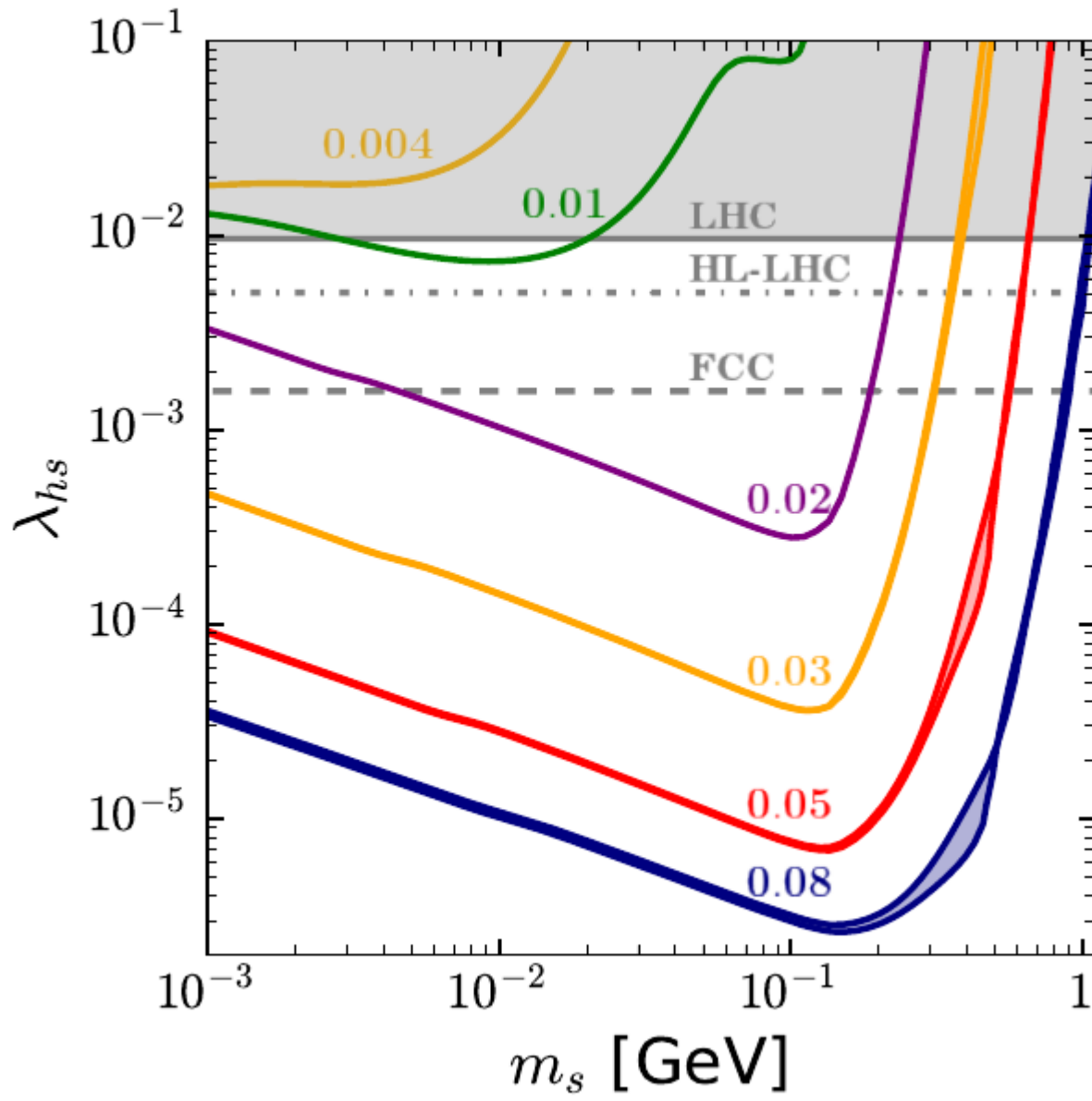


New input: treat **unknown** T_R as a free parameter \rightarrow direct detection + invisible Higgs decay

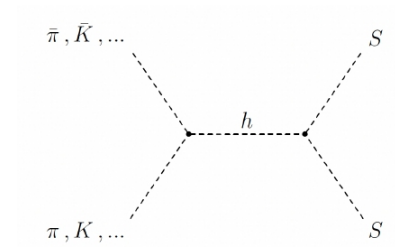
Very low T :



$$\Gamma_{\text{SM} \rightarrow \text{SS}} = \Gamma_{\text{SS} \rightarrow \text{SM}}^{\text{th}} = \frac{T}{2^5 \pi^4 m_h^4} \int_{4m_s^2}^{\infty} ds \sqrt{s(s - 4m_s^2)} K_1(\sqrt{s}/T) \Gamma_h(m_h = \sqrt{s}) |\mathcal{M}_{\text{SS} \rightarrow h}|^2$$



***pion-dominated
production***

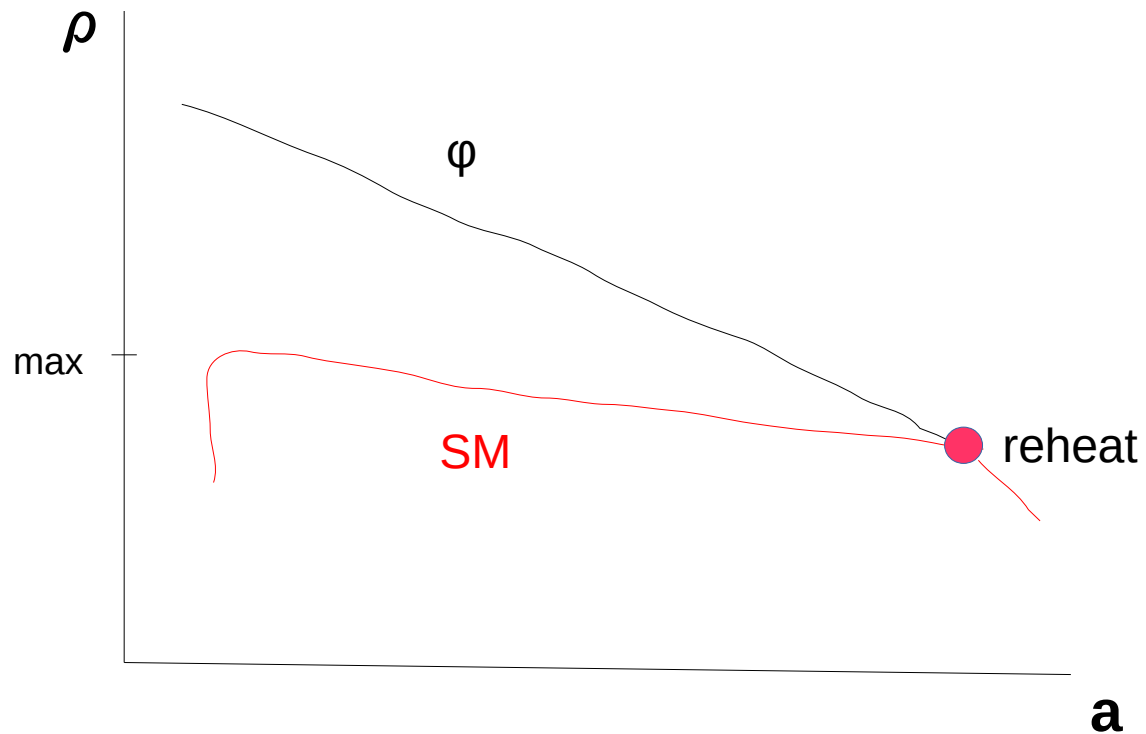


Detour: Reheating vs Maximal temperature

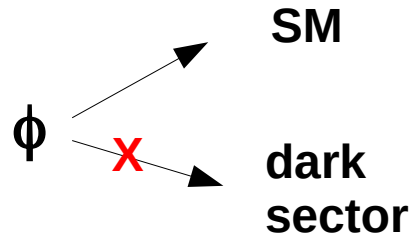
Not much is known if any ...

Textbook ($\phi \rightarrow SM$) :

$$T_{\max}^2 \sim M_{\text{pl}} (\Gamma_{\phi} H_0)^{1/2}$$

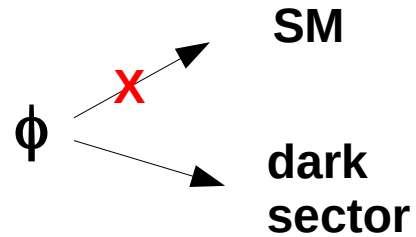


Closer look at “freeze-in” paradigm:



2 tiny couplings: DM-SM and DM-inflaton !

Logical extension:



T_{\max} changes completely!

SM ρ produced via “spectator” χ decay:

$$\dot{\rho} + 4H\rho = \Gamma_{\chi}\rho_{\chi}$$

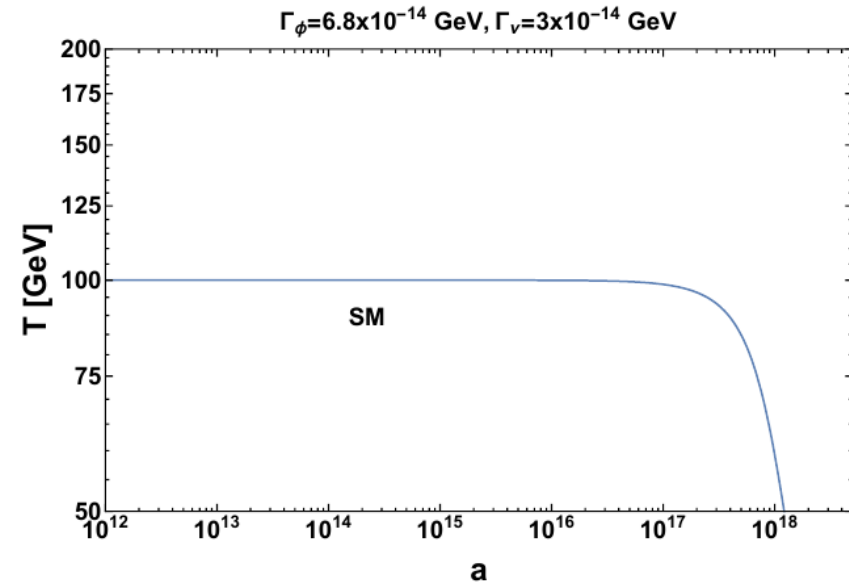
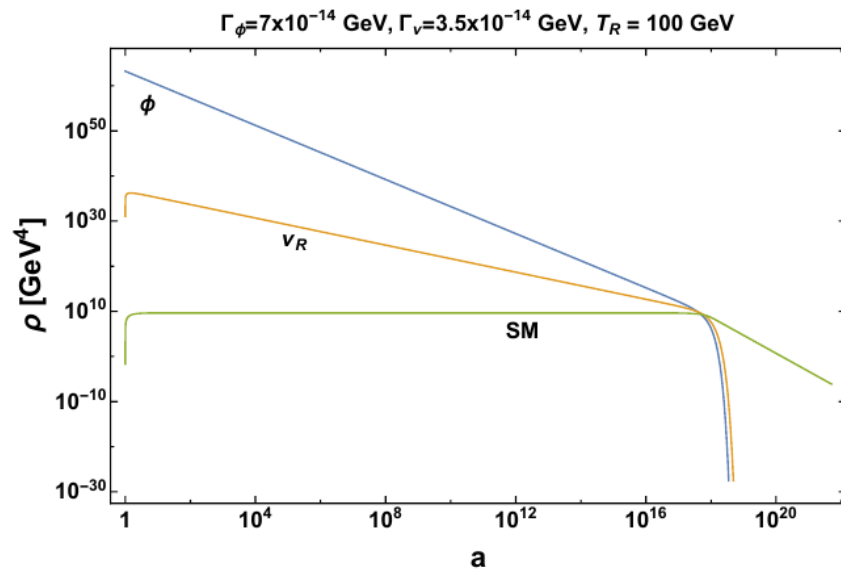
$$H = H_0/a^m,$$

$$\rho_{\chi} = \rho_{\chi}^0/a^n,$$

$$\rho(a) = \frac{\Gamma_{\chi}\rho_{\chi}^0}{(4-n+m)H_0} \left[\frac{1}{a^{n-m}} - \frac{1}{a^4} \right] \rightarrow \frac{\Gamma_{\chi}\rho_{\chi}^0}{(4-n+m)H_0} \frac{1}{a^{n-m}}$$

Example:

$$\phi \rightarrow \nu_R \nu_R, \quad \nu_R \rightarrow \text{SM}$$



$$T_R \sim T_{\text{max}}$$

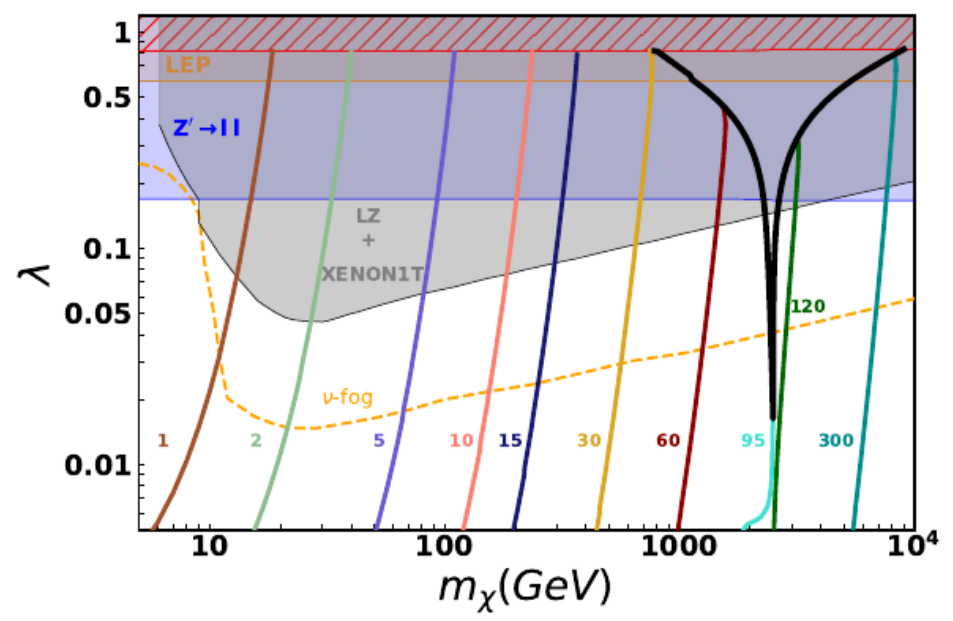
Z' extension

$$\mathcal{L} \supset -m_\chi \bar{\chi} \chi - \frac{1}{2} M_{Z'}^2 Z'_\mu Z'^\mu + \bar{\chi} \gamma^\mu (V_\chi - A_\chi \gamma_5) \chi Z'_\mu + \sum_f \bar{f} \gamma^\mu (V_f - A_f \gamma_5) f Z'_\mu$$

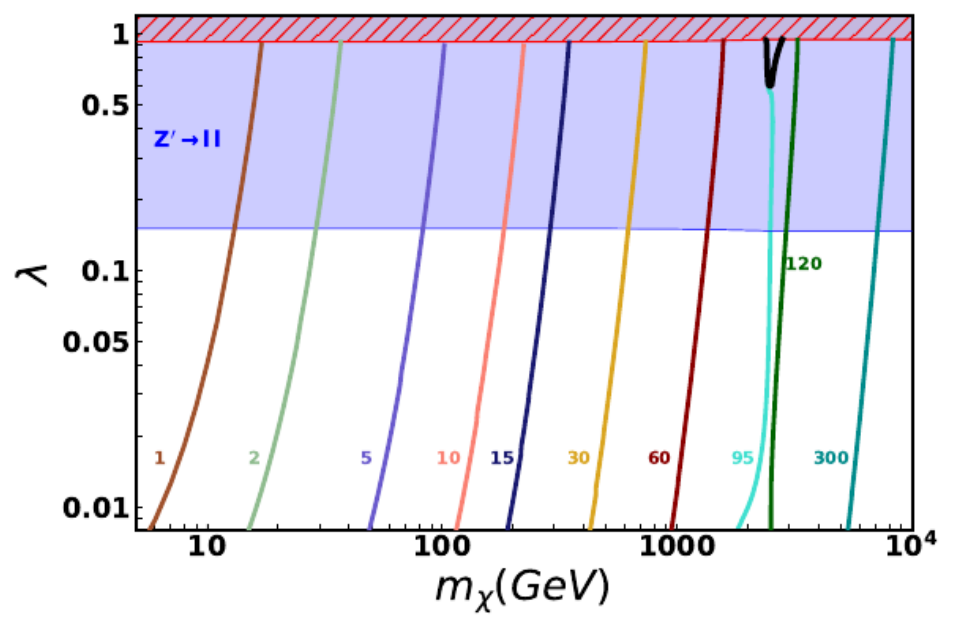
DM

Example: universal coupling λ , $M_{Z'} = 5$ TeV

vector



axial

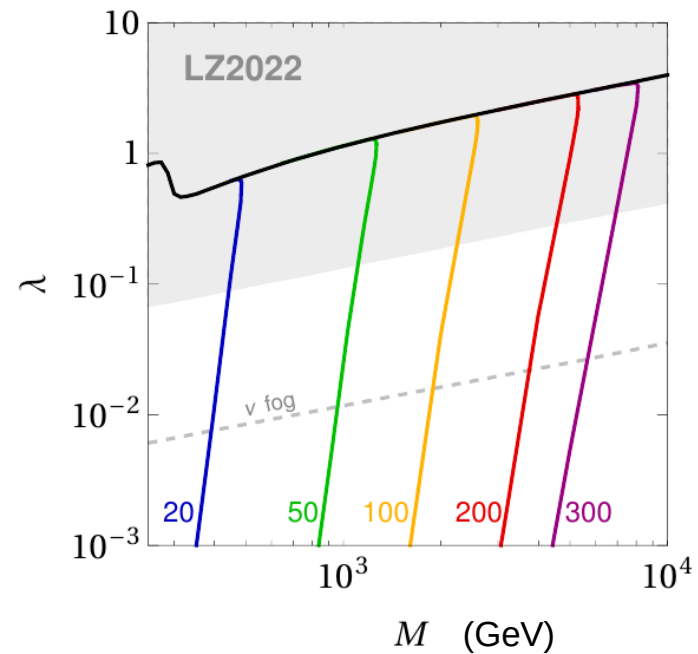
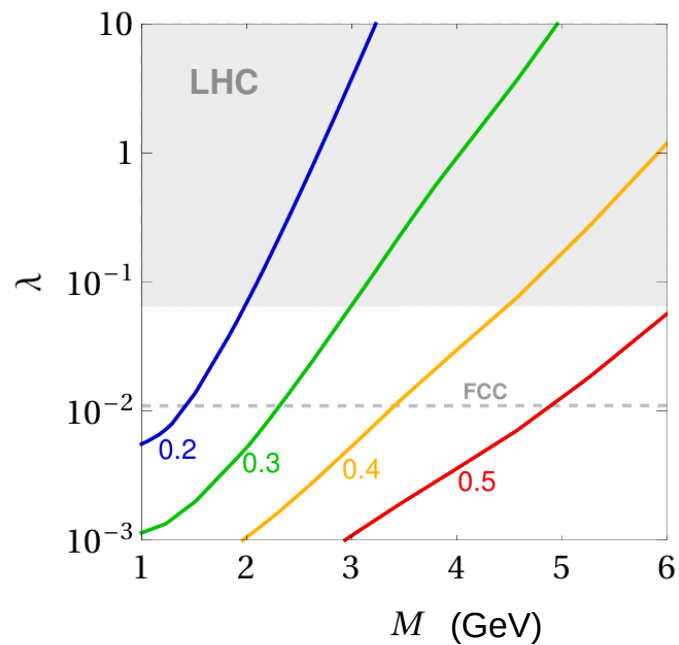


ν_R extension

$$\mathcal{L} \supset \frac{1}{2} M \bar{\nu} \nu + \frac{1}{2} \lambda S \bar{\nu} \nu + \left(y_{\nu,i} \bar{L}_{L,i} \tilde{H} \nu + \text{h.c.} \right)$$

DM

Example: scalar mixing = 0.2, heavy scalar mass = 300 GeV



CONCLUSION

- *dark relics are (over)produced during/after inflation*
- *non-thermal DM is sensitive to (quantum) gravity*
- *motivates freeze-in at stronger coupling (LHC, DD)*

