

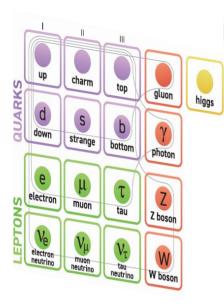
Naturalness-guided Search for the Origins of Matter

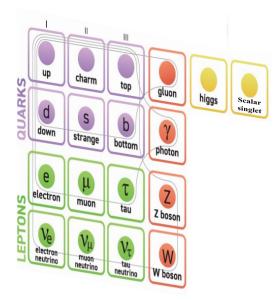
Ignacy Nałęcz



University of Warsaw Seminar on Particle Physics and Cosmology 06/06/2024, Warsaw

Based on: JHEP 2023.2:1-36 with M. Badziak





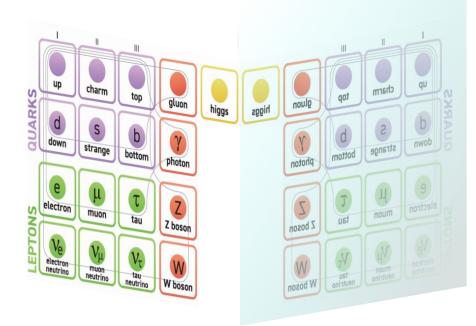


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- 2. First-order phase transition in Twin Higgs.
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- 6. Summary.

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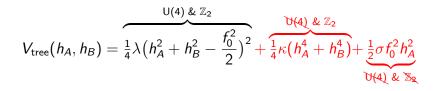
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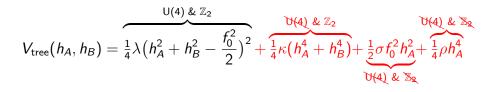
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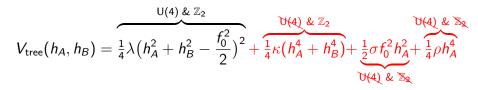
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$$V_{\text{tree}}(h_A, h_B) = \underbrace{\frac{U(4) \& \mathbb{Z}_2}{\frac{1}{4}\lambda (h_A^2 + h_B^2 - \frac{f_0^2}{2})^2}}_{U(4) \& \mathbb{Z}_2}$$

$$V_{\text{tree}}(h_A, h_B) = \underbrace{\frac{1}{4}\lambda(h_A^2 + h_B^2 - \frac{f_0^2}{2})^2}_{\frac{1}{4}\kappa(h_A^4 + h_B^4)} + \underbrace{\frac{1}{4}\kappa(h_A^4 + h_B^4)}_{\frac{1}{4}\kappa(h_A^4 + h_B^4)}$$







• f_0 is constrained by the LHC data and naturalness $3v_{SM} < f_0 < 8v_{SM}$.

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• Higgs mechanism requires U(4) to be the approximate symmetry $|\kappa|, |\sigma|, |\rho| < \lambda$.

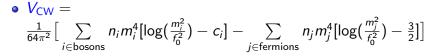
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$V_{\text{eff}}(h_A, h_B, T) = V_{\text{tree}}(h_A, h_B) + V_{\text{CW}}(h_A, h_B) + V_{\text{therm}}(h_A, h_B, T),$

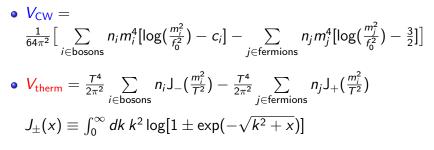
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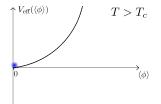


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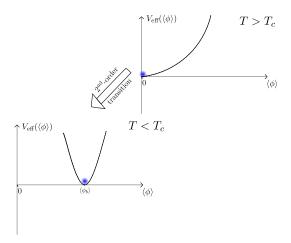
where

• $V_{CW} = \frac{1}{64\pi^2} \Big[\sum_{i \in \text{bosons}} n_i m_i^4 \Big[\log(\frac{m_i^2}{f_0^2}) - c_i \Big] - \sum_{j \in \text{fermions}} n_j m_j^4 \Big[\log(\frac{m_j^2}{f_0^2}) - \frac{3}{2} \Big] \Big]$ • $V_{\text{therm}} = \frac{T^4}{2\pi^2} \sum_{i \in \text{bosons}} n_i J_-(\frac{m_i^2}{T^2}) - \frac{T^4}{2\pi^2} \sum_{j \in \text{fermions}} n_j J_+(\frac{m_i^2}{T^2})$ $J_{\pm}(x) \equiv \int_0^\infty dk \ k^2 \log[1 \pm \exp(-\sqrt{k^2 + x})]$ • Daisy diagrams resummation: $m_i^2 \to \bar{m}_i^2 = m_i^2 + \Pi^2(T)$ in V_{CW} and V_{therm} .

Phase transition types



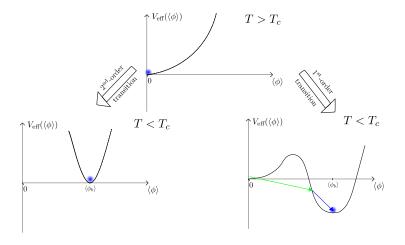
Phase transition types

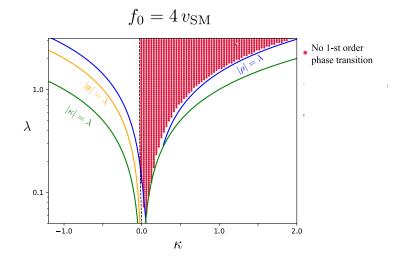


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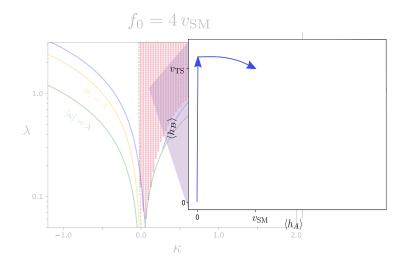
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Phase transition types

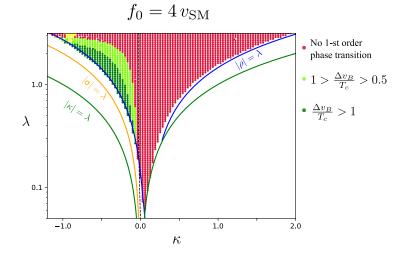


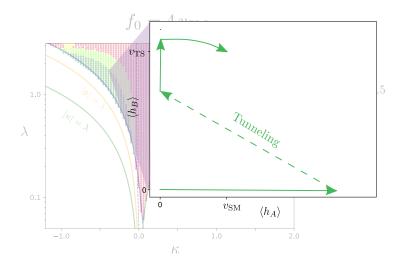


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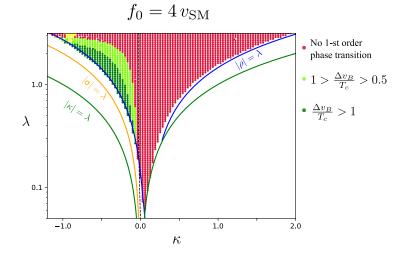


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Direction of the symmetry breaking

$$\begin{cases} 0 = \zeta_{\text{SM}} T_{\text{SM}}^2 - \lambda f_0^2 (1 - \frac{\sigma}{\lambda}) \\ 0 = \zeta_{\text{TS}} T_{\text{TS}}^2 - \lambda f_0^2 \end{cases}$$

[1] K. Fujikura et al. "Phase transitions in twin Higgs models", JHEP, 2018.12 (2018): 1-35.

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$\mathbb{Z}_2\text{-breaking}$ in twin lepton Yukawas

• Enhancment of the twin Yukawa couplings makes $\zeta_{\text{TS}} > \zeta_{\text{SM}}$ in $\sqrt{\frac{\zeta_{\text{TS}}}{\zeta_{\text{SM}}}(1-\frac{\sigma}{\lambda})}$ and thus allow for $T_{\text{SM}} > T_{\text{TS}}$, even for *positive* σ .

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We consider two variants:

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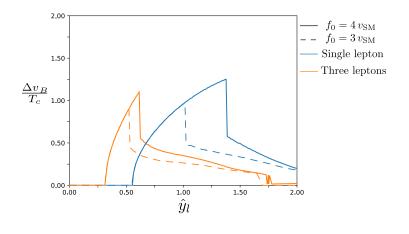
• The UV agnostic TH with enhanced lepton Yukawas $\hat{y}_l \gg y_l$. • The MSSM*-like extension with scalar lepton partners

$$\begin{split} \hat{m}_{\mathsf{sl}\ R}^2 &= \hat{\mu}_R^2 + \frac{1}{2}\tilde{y}_I^2h_B^2\cos^2\beta - \frac{1}{4}g'^2h_B^2\cos(2\beta), \\ \hat{m}_{\mathsf{sl}\ L}^2 &= \hat{\mu}_L^2 + \frac{1}{2}\tilde{y}_I^2h_B^2\cos^2\beta - \frac{1}{8}(g^2 - g'^2)h_B^2\cos(2\beta). \end{split}$$

*MSSM-Minimal Supesymmetric Standsrd Model

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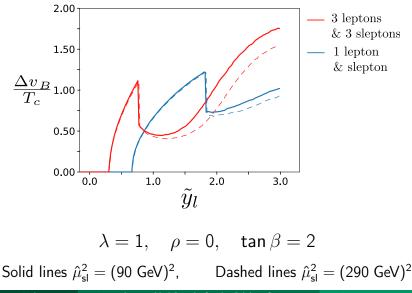
UV agnostic TH with enhanced twin lepton Yukawas



 $\lambda = 1, \quad \rho = 0$

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MSSM-like TH



Ingredients (aka Sakharov Conditions)

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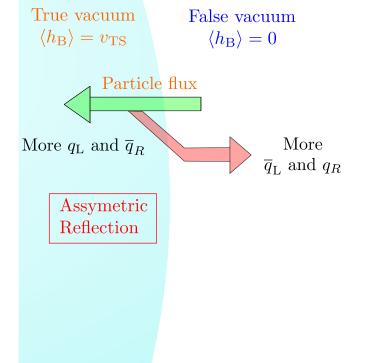
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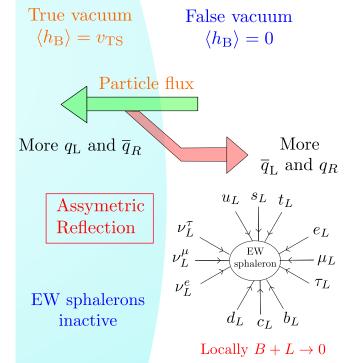
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True vacuum $\langle h_{\rm B} \rangle = v_{\rm TS}$

False vacuum $\langle h_{\rm B} \rangle = 0$

 $\vec{\xi_w}$





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Necessary ingredients

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- Transport operator there are many suitable options, for instance neutron operator

$$\mathcal{O}_n = \frac{1}{M^5} \bar{u}_R \bar{d}_R \bar{d}_R \hat{u}_L \hat{d}_L \hat{d}_L + \text{h.c.},$$

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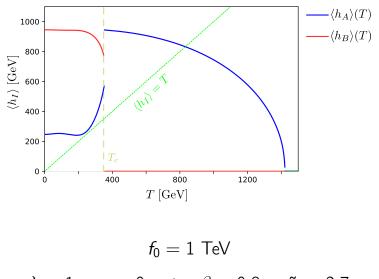
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Washout?

 $\hat{B} = X, \quad B = 0$ After FOPT ends $\hat{B} \approx \frac{X}{2}, \quad B \approx \frac{X}{2}$ \mathcal{O}_n in equilibrium \mathcal{O}_n decoupled, SM sphaleros active $B + L = 0, B - L \approx B \approx \frac{X}{4}$ More accurate computation in equilibrium approximation yields $B \approx 0.24X$. June 6, 2024 15 / 26

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At $T \ll f_0$ the effective field theory approximation fixes h_B

$$h_B pprox \sqrt{f_0^2 - h_A^2}.$$

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Condition for the symmetry non-restoration in TH

$$\sum_{\in \text{fermions}} \frac{n_j}{4} (\hat{y}_j^2 - y_j^2) \ge 5,$$

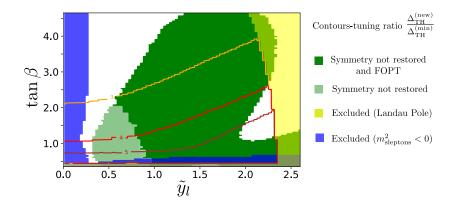
[2] O. Matsedonskyi, "High-Temperature Electroweak Symmetry Breaking by SM Twins", JHEP, 10.1007 (2021): 4-36.

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Twin leptons, quarks and scalar partners

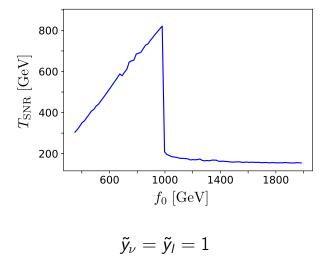


$$\lambda = 1, \ \rho = 0,$$

$$\tilde{y}_{\nu} = 0.9, \ \hat{y}_{u} = \hat{y}_{c} = 0.4, \ \hat{y}_{s} = \hat{y}_{b} = 0.3$$

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Temperature of EW symmetry breaking



 $\lambda = 1, \ \rho = 0, \quad \hat{y}_u = \hat{y}_c = 0.4, \quad \hat{y}_s = \hat{y}_b = 0.3$

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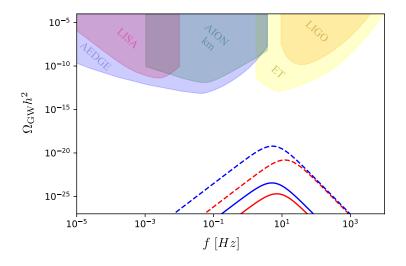
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- The new generation of GW detectors will partially probe the range at which gravitational signal from scalar FOPTs is expected.

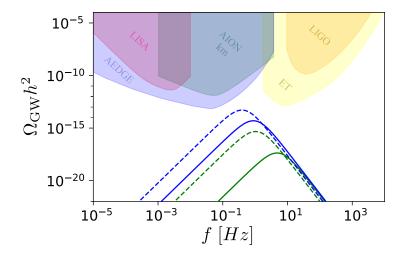
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Gravitational waves in UV agnostic TH



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Gravitational waves in SUSY TH



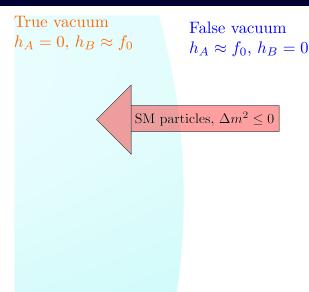
Non-equilibrium effects

True vacuum $h_A = 0, \ h_B \approx f_0$

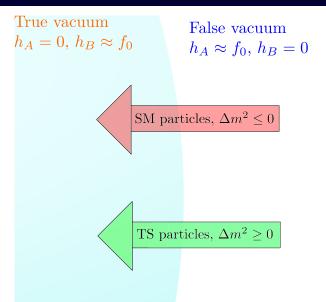
False vacuum $h_A \approx f_0, h_B = 0$

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Non-equilibrium effects



Non-equilibrium effects



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- In supersymmetric TH extensions the emission of detectable gravitational signal is not excluded.