

Talk at CGPG Feb 2001

Isolated Horizons and Their Secrets.

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- Geometry and invariants of non-expanding horizons
 - invariant description of data evolving tangentially,
 - invariant frame,
 - crossover spheres of WIHs and the crossover sphere

- Isolated Horizons
 - When a non-expanding horizon is isolated?
 - The uniqueness issue -special IHs of 2-dimensional symmetry group
 - the Kerr IH: invariant characterizations

- Neighborhood of IH
 - invariants,
 - Killing vector fields: the existence and nonexistence conditions

- Open Problems

- Sufficient conditions for the existence of Killing vector fields

- the degenerate cases,

- absorbing horizons.

Geometry and invariants of non-expanding horizons

Definitions and assumptions.

Definition: *Non-expanding horizon Δ is a null 3-submanifold of M such that:*

i) Δ is diffeomorphic to the product

$$\hat{\Delta} \times \mathbf{R}, \quad (1)$$

where $\hat{\Delta}$ is a 2-sphere, and the fibers of the projection

$$\hat{\Delta} \times \mathbf{R} \rightarrow \hat{\Delta} \quad (2)$$

correspond to the null geodesics in Δ ;

ii) the family of null geodesics tangent to Δ is non-expanding.

The energy assumptions and consequences. Assumptions:

$$T_{\mu\nu}\ell^\mu\ell^\nu \geq 0, \quad T_{\mu\nu}\ell^\nu \text{ causal} \quad (3)$$

$$\text{for } \ell, X \in T(\Delta), \quad \text{s.t. } \ell^\mu\ell_\mu = 0. \quad (4)$$

$$(5)$$

Consequences:

$$R_{a\nu}\ell^\nu X^a = 0, \quad X \in T(\Delta). \quad (6)$$

And ℓ is a double principal direction of the Weyl tensor.

Geometry of Δ . The induced metric q satisfies

$$\mathcal{L}_\ell q = 0 \quad (7)$$

It implies, that ∇ preserves the tangent bundle $T(\Delta)$, and defines a covariant derivative therein,

$$\mathcal{D}_X Y := \nabla_X Y. \quad (8)$$

Definition: The geometry of Δ is the pair (q, \mathcal{D}) .

\mathcal{D} is not determined by q , $[\mathcal{L}_\ell, \mathcal{D}] \neq 0$ in general.

Rotation of Δ , the geometry ingredient. There is a 1-form ω on Δ s. t.

$$\mathcal{D}\ell = \omega \otimes \ell \quad (9)$$

Definition: ω is called the rotation potential 1-form of (Δ, ℓ) , and the quantity

$$\kappa^{(\ell)} = \ell^a \omega_a \quad (10)$$

the surface gravity of ℓ .

$$\ell' = f\ell \Rightarrow \omega' = \omega + d\ln f \quad (11)$$

The invariant 2-form is,

$$d\omega = 2i\text{Im}\Psi_2 \text{vol}^{(2)}, \quad D\text{Im}\Psi_2 = 0. \quad (12)$$

can be considered as the angular velocity of Δ . (blackboard!) The zeroth law:

$$d\kappa^{(\ell)} \hat{=} \mathcal{L}_{\ell}\omega \quad (13)$$

This is a geometric version of the zeroth law of BH thermodynamics.

The remaining ingredients of \mathcal{D} . Deformations of the 2-sphere sections S_v of Δ as expanded in the

orthogonal-transversal null direction define the remaining components of \mathcal{D} , so called shear and expansion. Those depend on the pullback of Ricci tensor on S_v , and on initial conditions at any fixed cross-section of Δ .

Invariant choice of ℓ . For every geometry (q, \mathcal{D}) , there is ℓ , s.t.

$$d\kappa^{(\ell)} = \mathcal{L}_\ell \omega = 0, \quad (14)$$

but it is not unique. We call (Δ, ℓ) a Weakly IH. **BLACKBOARD Proposition:** For a given value of the surface gravity

$$\kappa^{(\ell)} = \kappa_0 \neq 0, \quad (15)$$

every generic non-expanding horizon (Δ, q, \mathcal{D}) such that

$$DR_{m\bar{m}} = 0 \quad (16)$$

admits a unique null, non-trivial WIH structure (Δ, ℓ) s. t. for every section S_v of Δ it is true

that the transport of S_ν by the flow of ℓ does not change the corresponding transversal expansion μ . The genericity condition is that 0 is not an eigenvalue of the following operator $M : L^2(S) \rightarrow L^0(S)$,

$$M := \hat{\Delta} + 2\hat{\omega}^A \partial_A + \hat{\text{div}}\hat{\omega} + \hat{\omega}_A \hat{\omega}_B \hat{q}^{AB} - K + R_{m\bar{m}}, \quad (17)$$

and the non-triviality is that ℓ can not be identically 0 on a null geodesic.

Geometric version of the proposition: crossover sphere

Consider the analytic-geodesic completion of Δ and (q, \mathcal{D}) . For every ℓ such that

$$\kappa^{(\ell)} = \kappa_0 \neq 0, \quad (18)$$

The equation

$$\ell = 0, \quad (19)$$

defines a crossection S_ℓ of the completion.

Proposition: A null vector field tangent to Δ is the invariant vector field ℓ of the previous proposition if

and only if $\kappa^{(\ell)} = \kappa_0$ and the zero set of ℓ has zero expansion.

We will call such cross-section the crossover sphere of Δ .

This solves the issue of invariant evolution of the data: it is

$$[\mathcal{L}_\ell, \mathcal{D}]^a{}_{bc} = \ell^a(\dot{\lambda}m_b m_c + c.c.). \quad (20)$$

Good cuts foliation. Given (q, \mathcal{D}) consider a space-like foliation of Δ preserved by the flow of the invariant vector field ℓ . The pullback of the rotation 1-form potential ω to the slices, defines a unique 1-form $\hat{\omega}$ on S ,

$$\hat{\omega} = \hat{*}dU + \hat{d}p. \quad (21)$$

The first term is foliation invariant, the second is arbitrary.

Definition We call a foliation the good cuts foliation if

$$\hat{\omega} = \hat{*}dU. \quad (22)$$

For every ℓ such that $\kappa^{(\ell)} = \text{const}$, there is a unique good cuts foliation.

Invariant frame on Δ . The invariant ℓ and corresponding good cuts foliation uniquely define a transversal, null vector field n , such that

$$n_{\mu}\ell^{\mu} = -1, \quad n \text{ orthogonal to the foliation.} \quad (23)$$

Finally, we can complete (n, ℓ) to a null frame (m, \bar{m}, n, ℓ) by using the Gauss curvature K of the crosssections, wherever $dK \neq 0$,

$$(\delta - \bar{\delta})K = 0. \quad (24)$$

Proposition The above invariant null frame is uniquely and naturally defined by the geometry (q, \mathcal{D}) . (diffeos: blackboard!)

Isolated Horizons.

Definition A non-expanding horizon (Δ, q, \mathcal{D}) and a null, tangent vector field ℓ are called IH if

$$[\mathcal{L}_\ell, \mathcal{D}] = 0. \quad (25)$$

The existence Issue:

(i) A non-expanding horizon (Δ, q, \mathcal{D}) admits the IH structure if and only if the invariant vector field ℓ defines it.

(ii) We have also derived the explicit existence conditions.

True degrees of freedom In the $\kappa^{(\ell)} \neq 0$ case, the geometry (q, \mathcal{D}) is explicitly given by: Gauss curvature K , the rotation potential U , the Ricci tensor components given by contraction with the vectors tangent to Δ .

The crossover sphere Proposition: If $(\Delta, [\ell])$ is an isolated horizon, then the crossover sphere of the completion of Δ has zero shear and expansion in any orthogonal null direction. Conversely, if the completion of a non-expanding and shear free horizon Δ contains a cross-section that is shear free and non-expanding in every orthogonal null direction, then, the corresponding null vector field ℓ is an isolated horizon.

The uniqueness issue Given an isolated horizon (Δ, q, \mathcal{D}) , are there ℓ and $\ell' = f\ell$ ($df \neq 0$) such that

$$[\mathcal{L}_\ell, \mathcal{D}] = 0 = [\mathcal{L}_{\ell'}, \mathcal{D}].? \quad (26)$$

Results: (i) The general case: we derived necessary non-uniqueness conditions (a non-uniqueness test).

(ii) Special cases in which the uniqueness has been proven:

- (a) *If the geometry of an isolated horizon is sufficiently close to that of the Kerr-Newman horizon*
- (b) *If an isolated horizon Δ is non-rotating and the Ricci tensor vanishes on Δ*
- (c) *If the Ricci tensor vanishes on an isolated horizon Δ , and the geometry (q, \mathcal{D}) is not the one characterized below.*

The isolated horizon of 2-dim null symmetry group.
The only non-unique IH structures are define on the following non-expanding horizon (Δ, q, \mathcal{D}) . Consider a vacuum isolated horizon (Δ, q, \mathcal{D}) and its null symmetry $\ell = \ell_0$ such that

$$\kappa^{(\ell)} = 0 \quad (27)$$

and it admites a crossection of the zero expansion and shear in each orthogonal null direction. This horizons geometry (q, \mathcal{D}) has 2- dimensional group of null symmetries.

*Proposition (Pawłowski):*In the cylindrically symmetric case, the only non-unique IH structures come from the modification of the external Kerr data.

Geometric characterization of the Kerr horizon.

*Proposition:*Each of the below conditions 1) and 2) is necessary and sufficient for an axially symmetric IH to be the Kerr IH:

1. Let Φ be the rotation Killing vector,

$$d(\Psi_2^{-\frac{1}{3}}) = a_0 \Phi \lrcorner \epsilon^{(2)} \quad (28)$$

2. $R_{\mu\nu}|_{\Delta} = \partial R_{\mu\nu}|_{\Delta}$ and the Weyl tensor is of the type D at Δ .

Neighborhood of IH

The unique extension of the good cuts foliation, invar

Given a horizon (non-expanding or isolated), the invariant null vector field ℓ , and the corresponding good cuts foliation we extend it in the following way: (blackboard)

(i) the good cuts foliation by the family of the orthogonal null geodesics

(ii) The vector field n , by

$$\nabla_n n = 0, \quad (29)$$

(iii) The vector field ℓ , to ξ , s.t

$$\xi|_{\Delta} = \ell, \quad \mathcal{L}_n \xi = 0. \quad (30)$$

The vector field ξ is not any longer null out of Δ .

Proposition If Δ is a Killing horizon, then ξ is the Killing vector field.

The unique extension of the null frame By

$$\nabla e^\mu = 0. \quad (31)$$

The Gauss coordinates A coordinates system (x^A, v) compatible with the good cuts foliation is naturally extended by

$$n^\mu \partial_\mu x^A = n^\mu \partial_\mu v = 0 \quad (32)$$

and completed by r ,

$$n^\mu \partial_\mu r = -1, \quad r|_\Delta = 0. \quad (33)$$

to coordinates (x^A, v, r) .

The invariants:

(i) The curvature of the 2-spheres $r = \text{const}$, $v = \text{const}$

(ii) the products $\xi^\mu \ell_\mu$, $\xi^\mu n_\mu$, $\xi^\mu m_\mu$

(iii) Any geometric tensor components.

The degrees of freedom: Friedrich's reduced data.

The degrees of freedom (q, U) in the geometry (q, \mathcal{D}) are compatible with the Friedrich-Rendall reduced data and provide true degrees of freedom for a neighborhood of Δ . In the $\kappa^{(\ell)} \neq 0$, vacuum case, they are:

(i) Geometry (q, \mathcal{D}) of Δ freely parametrized by q, U

(ii) R_{nmnm} on two transversal half-leaves of the invariant null foliation.

Symmetric horizons and the existence of the Killing vector fields conditions.

Symmetries of Δ .

Definition A vector field $k \in T(\Delta)$ generates a symmetry of Δ , if

$$\mathcal{L}_k q = 0, \quad L_k \mathcal{D} = 0.$$

Proposition The invariant null vector field ℓ , the corresponding good cuts foliation, and the ingredients we use to describe (q, \mathcal{D}) are preserved by the symmetries of Δ .

Every symmetry generating $k \in T(\Delta)$ is given by a Killing vector field of a single slice, and a constant a_0 ,

$$k = a_0 \ell + Y, \quad \mathcal{L}_\ell Y = 0, \quad \mathcal{L}_Y q = 0.$$

It follows from the evolution of \mathcal{D} along the null generators, that:

Proposition Suppose $k = a_0 \ell + Y$; then both the vector fields $a_0 \ell$ and Y generate symmetries of Δ .
(blackboard)

Classification of the symmetries of Δ :

Class 1. $k = \ell$. Isolated horizon

Class 2. $k = Y$. Axially symmetric non-expanding horizon.

Class 3. $k_1 = \ell$, $k_2 = Y$, Axially symmetric isolated horizon

Class 4. $k_1 = \ell$ and 3-dim rotation symmetry group.

Killing vector fields: possible candidates.

Proposition Suppose K is a Killing vector field tangent to Δ and defined in a neighborhood of Δ . Then:

1. K preserves the invariant foliation of the neighborhood, the vector field ξ and the invariant null frame.

2. K is uniquely determined by a symmetry k of Δ it generates, by

$$\mathcal{L}_n K = 0, \quad K|_{\Delta} = k.$$

Therefore, given a non-expanding horizon Δ , the natural null vector field ℓ , the good cuts foliation and the corresponding Gauss coordinates, the only possible candidates for a Killing vector field are,

$$K = \begin{cases} a_0 \partial_v = \xi, \\ b_0 \partial_\varphi, \\ a_0 \partial_v + b_0 \partial_\varphi, \end{cases} \quad (34)$$

the 2-sphere $\hat{\Delta}$ of the null generators coordinates $(x^A) = (\theta, \varphi)$ are adjusted to a Killing vector field of $\hat{\Delta}$.

Killing vector fields: the existence conditions, the vacu

1. Axial symmetry:

(i) Necessary condition at Δ : the axial symmetry of (q, \mathcal{D})

(ii) *Necessary and sufficient: the axial symmetry of the reduced data. trivial*

2. *The Killing horizon case:*

(i) *Necessary conditions at Δ : the functions $\partial_r^N R_{nmnm}$, $N = 0, 1, 2, 3, 4, \dots$ at a fixed cross-section of Δ are determined by (q, \mathcal{D}) in an explicit way.*

(ii) *Necessary and sufficient: the following equation to be satisfied by $\Psi_4 := \overline{R_{nmnm}}$ on a fixed leaf of the invariant foliation,*

$$(X^A \partial_A + H \partial_r) \Psi_4 = (\rho - 4\epsilon) \Psi_4 + (\bar{\delta} + 4\pi + 2\alpha) \Psi_3 - 3\lambda \Psi_2.$$

3. *The cross diagonal case ($K = a_0 \ell + \partial_\varphi$)*

The necessary conditions: both, the necessary conditions 1. and 2. above.

Since R_{mnmn} is free, in this way we control the non-expanding (isolated) horizons which do NOT admit any Killing vector field.