

# Two-Higgs Doublet Models with Scalar Singlet Dark Matter

Bohdan Grzadkowski  
University of Warsaw

"Summer School and Workshop on the Standard Model and Beyond",  
Corfu, September 9th 2014

## Outline:

- 2HDMS Model
- Motivations
- Strategy
- Resulting Constraints on the parameter space
- Direct DM detection constraints
- New Higgs physics at the LHC?
- Summary

2HDM: B. Dumont, J. Gunion, S. Kraml, Y. Jiang, arXiv:1405.3584

2HDMS: A. Drozd, B. Grzadkowski, J. F. Gunion and Y. Jiang, "Extending two-Higgs-doublet models by a singlet scalar field - the Case for Dark Matter", arXiv:1408.2106.

# 2HDM<sub>S</sub> model

## 2HDM<sub>S</sub> - Yukawa Interactions

- Type I (only  $H_2$  couples to fermions)
- Type II ( $H_2$  couples to up-type fermions,  $H_1$  other)

Symmetry:  $Z_2 : H_1 \rightarrow -H_1$ , other scalar fields  $Z_2$ -even  
 $Z'_2 : S \rightarrow -S$ , other fields  $Z'_2$ -even

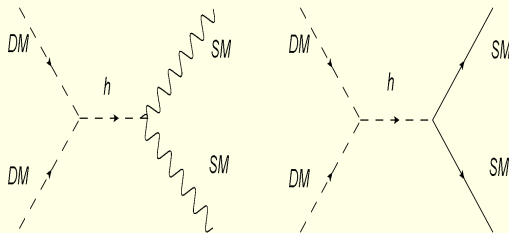
$$\begin{aligned} \mathcal{V} = & m_{11}^2 H_1^\dagger H_1 + m_{22}^2 H_2^\dagger H_2 - [m_{12}^2 H_1^\dagger H_2 + \text{h.c.}] + \frac{\lambda_1}{2} (H_1^\dagger H_1)^2 + \frac{\lambda_2}{2} (H_2^\dagger H_2)^2 \\ & + \lambda_3 (H_1^\dagger H_1) (H_2^\dagger H_2) + \lambda_4 (H_1^\dagger H_2) (H_2^\dagger H_1) + \left\{ \frac{\lambda_5}{2} (H_1^\dagger H_2)^2 + \text{h.c.} \right\} \\ & + \frac{m_0^2}{2} S^2 + \frac{\lambda_S}{4!} S^4 + \kappa_1 S^2 (H_1^\dagger H_1) + \kappa_2 S^2 (H_2^\dagger H_2) \end{aligned}$$

EWSB:  $Z'_2$  unbroken  $\rightarrow$  **NO VEV FOR S**

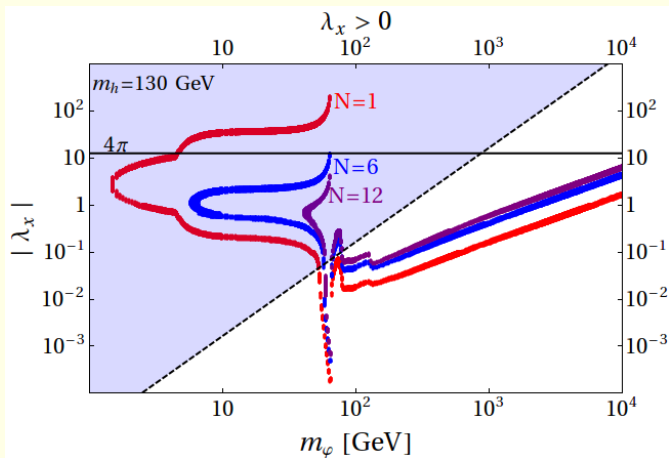
$$H_{1,2} = \begin{pmatrix} \varphi_{1,2}^+ \\ (v_{1,2} + \rho_{1,2} + i\eta_{1,2})/\sqrt{2} \end{pmatrix} \quad \tan \beta \equiv \frac{v_2}{v_1}, \quad v_1^2 + v_2^2 = (246 \text{ GeV})^2$$

## 2HDMS

- An attempt to provide both extra CP violation *and* DM candidate - 2HDMS minimal model,
- 2HDM provides an interesting "low-mass" new physics accessible at the LHC,
- To have a chance for  $M_{DM} < m_h/2$



# Motivations



$$BR(h \rightarrow \varphi\varphi) \propto \lambda_x^2 \quad \text{for} \quad V(H, \varphi) = \dots + \lambda_x H^\dagger H \varphi^2$$

5 mass eigenstates:  $h, H, A, H^\pm, S$

- 10 parameters in the potential, various basis possible

## General Basis:

- $\lambda_1, \lambda_2, \lambda_3, \lambda_4, \lambda_5$
- $m_{12}^2, \tan \beta$
- $m_S, \kappa_1, \kappa_2$

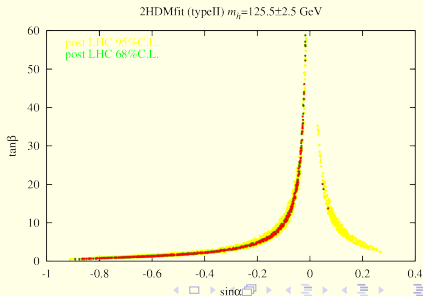
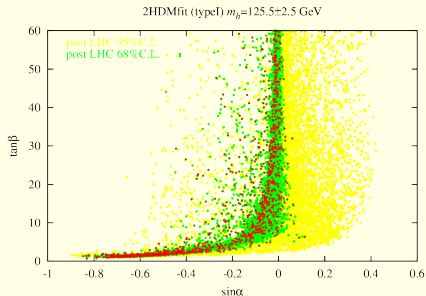
## Physical Basis:

- $m_h, m_H, m_A, m_{H^\pm}, \sin \alpha$
- $m_{12}^2, \tan \beta$
- $m_S, \lambda_h, \lambda_H$

- 2 types of Yukawa interaction

## 2HDM: Dumont, Gunion, Jiang, Kraml

- theoretical constraints  
(perturbativity, vacuum stability, perturbative unitarity)
- experimental constraints
  - B/LEP limits  $H^+$
  - STU
  - heavy Higgs search
  - LHC fit at 68% CL



## 2HDM

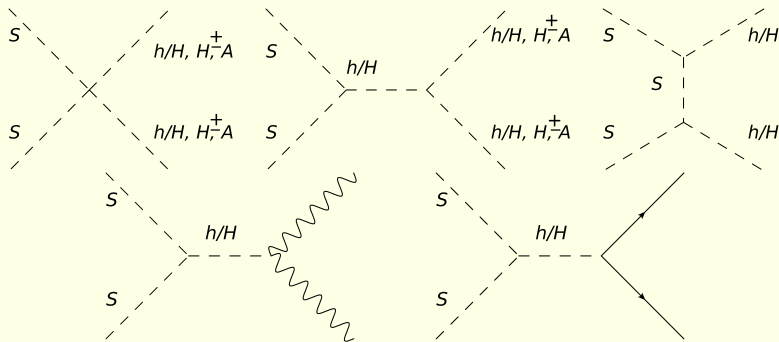
Take good 2HDM points

### Scalar Singlet parameter scan:

- $m_S \in [1 \text{ GeV}, 1 \text{ TeV}]$
- $\lambda_h, \lambda_H \in [-4\pi, 4\pi]$
- theoretical constraints (perturbativity, vacuum stability, perturbative unitarity, EWSB)
- with  $BR(h \rightarrow DM, DM) < 10\%$
- WMAP/Planck
- direct DM detection



# Strategy

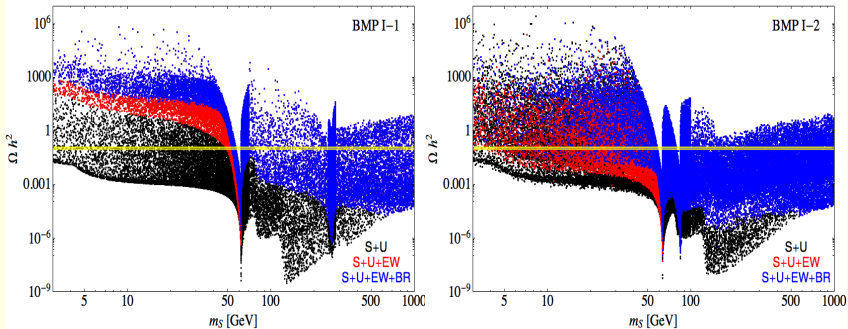


Calculation of DM relic abundance  $\Omega$ :

MicrOmegas by G. Belanger, F. Boudjema, A. Pukhov, A. Semenov,  
arXiv:0803.2360

$$\Omega^{WMAP/Planck} = 0.1187 \pm 0.0017$$

# Resulting Constraints on the parameter space

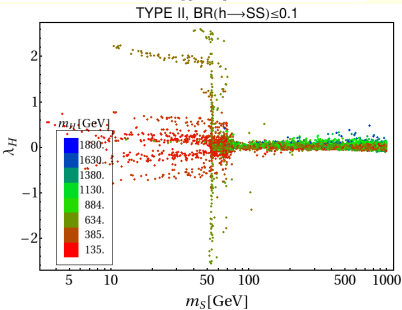
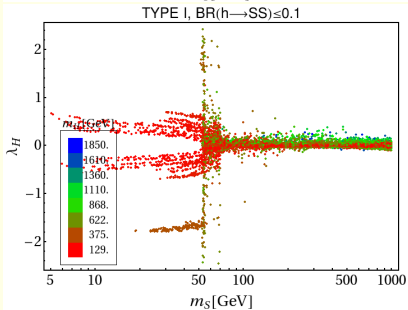
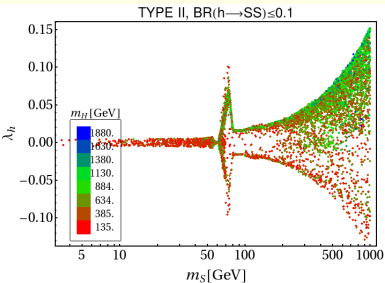
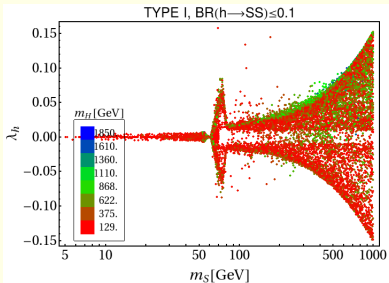


#	$\tan \beta$	$\sin \alpha$	$m_{12}^2$	$m_h$	$m_H$	$m_A$	$m_{H\pm}$
I-1	1.586	-0.587	+5621	123.71	534.25	645.13	549.25
I-2	1.346	-0.663	-2236	126.49	168.01	560.92	556.94

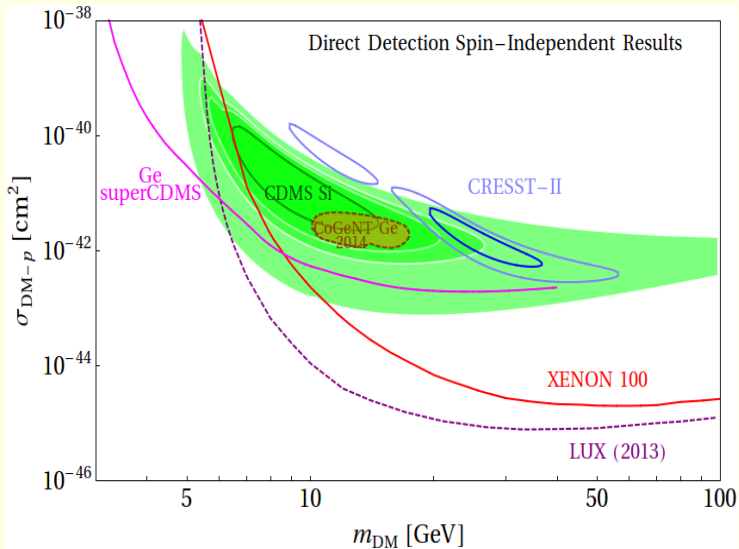
$$BR(h \rightarrow SS) = ???$$

- $\Omega_{DM}$  requires sufficiently strong SM - DM coupling
- search  $\lambda_h, \lambda_H$  give appropriate  $BR(h \rightarrow SS)$  i  $\Omega_{DM}$
- H responsible for DM production!

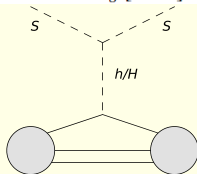
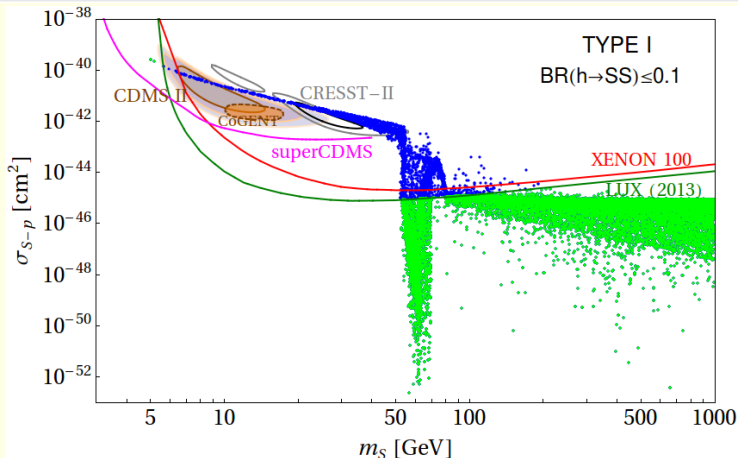
# Resulting Constraints on the parameter space



# Direct DM detection constraints



# Direct DM detection constraints



# Direct DM detection constraints

## TYPE II

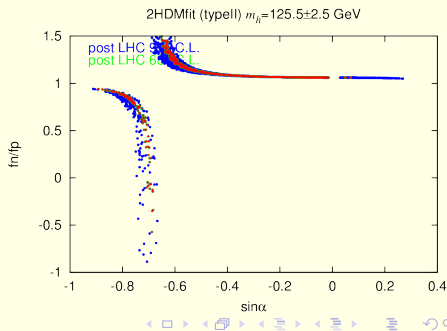
$$\sigma_{DM-N} = \frac{4\mu_{ZA}^2}{\pi} f_p^2 \left[ Z + \frac{f_n}{f_p} (A - Z) \right]^2$$

$$BR(h \rightarrow SS) \leq 0.1 \Rightarrow \lambda_h < 0.015$$

$$\frac{f_n}{f_p} = \frac{m_n}{m_p} \frac{\sum_q \left[ \left( \frac{\lambda_h}{\lambda_H} \xi_h^q + \left( \frac{m_h}{m_H} \right)^2 \xi_H^q \right) f_n^q \right]}{\sum_q \left[ \left( \frac{\lambda_h}{\lambda_H} \xi_h^q + \left( \frac{m_h}{m_H} \right)^2 \xi_H^q \right) f_p^q \right]} \rightarrow \frac{m_n}{m_p} \frac{\sum_q [(\xi_h^q + \xi_H^q) f_n^q]}{\sum_q [(\xi_h^q + \xi_H^q) f_p^q]} \quad (\text{S indep.})$$

**Table:** Yukawa couplings of up and down type quarks to light and heavy Higgs bosons  $h, H$  in Type I/II models. The Yukawa Lagrangian is normalised as follows:  $\mathcal{L}^{Yukawa} = \frac{m_q}{v} \xi_h^q \bar{q} q h + \frac{m_q}{v} \xi_H^q \bar{q} q H$

	Type I	Type II
$\xi_h^u$	$\cos \alpha / \sin \beta$	$\cos \alpha / \sin \beta$
$\xi_h^d$	$\cos \alpha / \sin \beta$	$-\sin \alpha / \cos \beta$
$\xi_H^u$	$\sin \alpha / \sin \beta$	$\sin \alpha / \sin \beta$
$\xi_H^d$	$\sin \alpha / \sin \beta$	$\cos \alpha / \cos \beta$



# Direct DM detection constraints

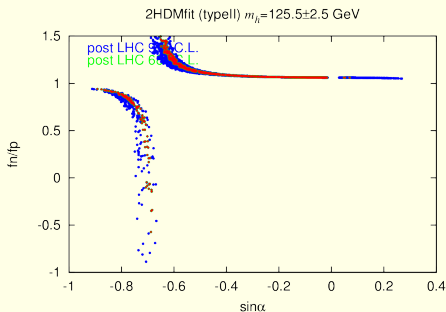
## TYPE II

$$\sigma_{DM-N} = \frac{4\mu_{ZA}^2}{\pi} f_p^2 \left[ Z + \frac{f_n}{f_p} (A - Z) \right]^2 \quad \sigma_{DM-p}^{EXP} \geq \sigma_{DM-p}^{THEO} \Theta^{EXP}(f_n, f_p)$$

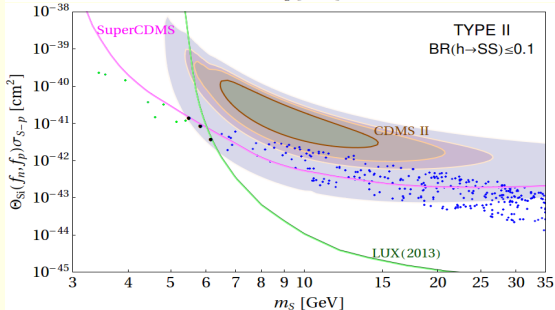
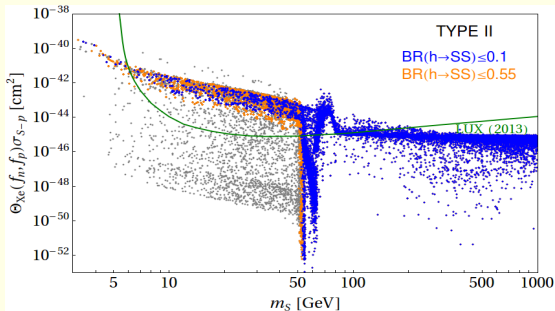
$$\Theta^{EXP}(f_n, f_p) = \sum_I \mu_I \left( \frac{Z_I}{A_I} + \frac{f_n}{f_p} \frac{A_I - Z_I}{A_I} \right)^2$$

**Table:** Yukawa couplings of up and down type quarks to light and heavy Higgs bosons  $h, H$  in Type I/II models. The Yukawa Lagrangian is normalised as follows:  $\mathcal{L}^{Yukawa} = \frac{m_q}{v} \xi_h^q \bar{q} q h + \frac{m_q}{v} \xi_H^q \bar{q} q H$

	Type I	Type II
$\xi_h^u$	$\cos \alpha / \sin \beta$	$\cos \alpha / \sin \beta$
$\xi_h^d$	$\cos \alpha / \sin \beta$	$-\sin \alpha / \cos \beta$
$\xi_H^u$	$\sin \alpha / \sin \beta$	$\sin \alpha / \sin \beta$
$\xi_H^d$	$\sin \alpha / \sin \beta$	$\cos \alpha / \cos \beta$



# Direct DM detection constraints





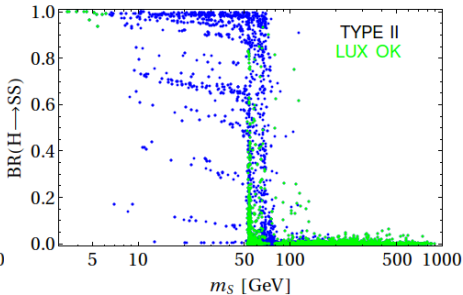
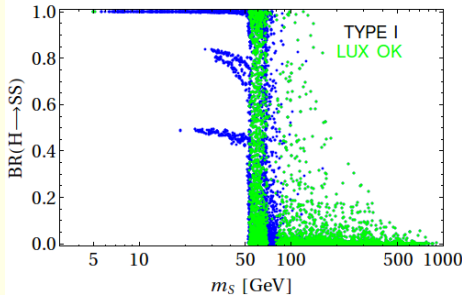
# Direct DM detection constraints

$\tan \beta$	$\sin \alpha$	$m_H$	$m_A$	$m_{H^\pm}$	$m_{12}^2$	$m_S$
2.092	-0.41	138	451	399	-12642	3.44; 3.56; 3.95
3.121	-0.282	187	546	571	8943	4.82; 5.48
2.192	-0.394	209	488	503	7518	5.40
1.728	-0.476	177	318	389	9382	5.16
1.789	-0.461	198	420	430	-6594	4.44; 5.15
1.488	-0.528	157	553	576	-10094	4.61
2.375	-0.363	259	260	339	15899	5.83

**Table:** Summary of the properties of the 2HDM Type II points which make it possible to realize  $m_S < 50$  GeV in agreement with within 99% CL for CDMS II imposing the full set of constraints including the LUX and SuperCDMS bounds and. All masses are given in GeV units.

# New Higgs physics at the LHC?

$H \rightarrow SS$  decay - invisible H!  
 $m_H \sim 130 - 200$  GeV



# Conclusions

- 2HDM is allowed by current collider limits, even in the non-decoupling regime
- 2HDMS provides a viable DM candidate and an opportunity for extra CP-violation
- 2HDMS is allowed by current collider and  $\Omega$  limits
- LUX requires  $m_S \gtrsim 50 \text{ GeV}$  (TYPE I, II) or together with SuperCDMS  $m_S \lesssim 6 \text{ GeV}$  (TYPE II)
- CDMS II requires  $|\lambda_h| < 0.05$ ,  $|\lambda_H| > 0.1$ , and implies large  $BR(H \rightarrow SS)$  (TYPE I, II)
- A fit of 2HDMS to LUX, superCDMS and CDMS II is only possible within 99% CL for CDMS II, for TYPE II model, then  $m_S \sim 3.4 - 5.8 \text{ GeV}$ . For those points  $BR(H \rightarrow SS) \gtrsim 90\%$